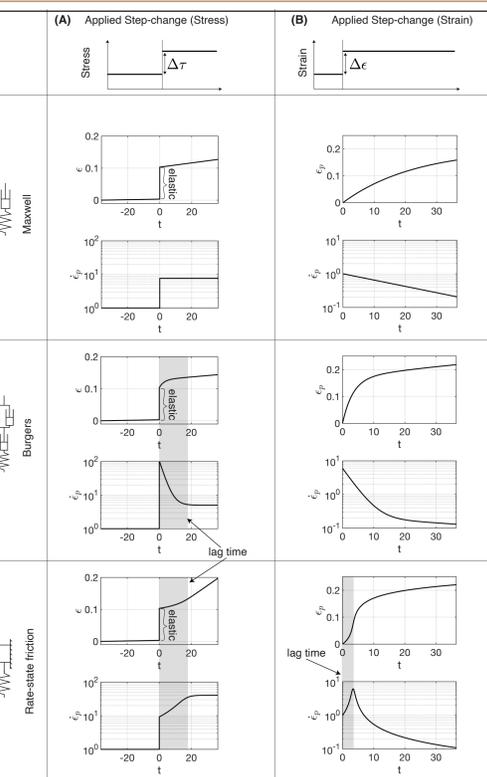


Main Points

1. We developed a numerical method to solve for the evolution of stress, strain rate and surface displacements in response to periodic and aperiodic earthquake sequences for popular linear and non-linear viscoelastic rheologies.
2. Homogeneous linear Maxwell materials are incompatible with typical interseismic and postseismic geodetic observations, while heterogeneous linear Maxwell, homogeneous linear Burgers and power-law materials are nearly indistinguishable for periodic cycles.
3. Earthquake sequence deformation (coseismic events spanning an order of magnitude) may be used to discriminate between linear and power-law rheologies, even with current uncertainties in geodetic observations.
4. Surface deformation kinematics may appear similar for different rheologies, but the details of regional stress transfer and rates of creep migration may change dramatically.



Earthquake cycles and rheological models (elasticity + ?)

The solid Earth behaves as a brittle-elastic material over short timescales (1 second - 1 day). At longer timescales, the viscous properties of the lithosphere-asthenosphere system control the mode of deformation. Together, the elastic and non-elastic (friction/viscous) properties of the solid Earth control the amplitude and timescales of stress transfer within the entire medium.

Improving estimates of these parameters has important implications for understanding fundamental issues such as the existence of plate-like tectonics, as well as more applied challenges in regional seismic hazard and risk.

- Observations that constrain the non-elastic properties of the Earth's lithosphere are limited.
- Postseismic relaxation of the lithosphere following large earthquakes allow us to use surface observations to constrain the non-elastic properties of the lithosphere-asthenosphere system.
- The postseismic relaxation problem is mathematically similar to a strain-step IVP (stress relaxation - Figure 1B). This makes it difficult to distinguish between various lab-derived rheological models (linear Maxwell, power-law, linear/power-law Burgers, rate-state friction) using surface deformation observations (Figure 1).
- We want to know - to what degree do current observations allow us to constrain an average unique rheological model of the lithosphere (and relevant parameters)?

Figure 1. Comparison of the response of three popular rheological descriptions (Maxwell, Burgers and rate-state friction) to a step-change in (A) stress and (B) strain. ϵ - total strain, ϵ_p - non-elastic strain. The stress-step experiments allow visual discrimination between the various rheologies, while strain-step experiments mostly appear as relaxation curves with monotonically decaying rates (except for rate-state friction).

Numerical Model (Viscoelastic earthquake sequence simulations)

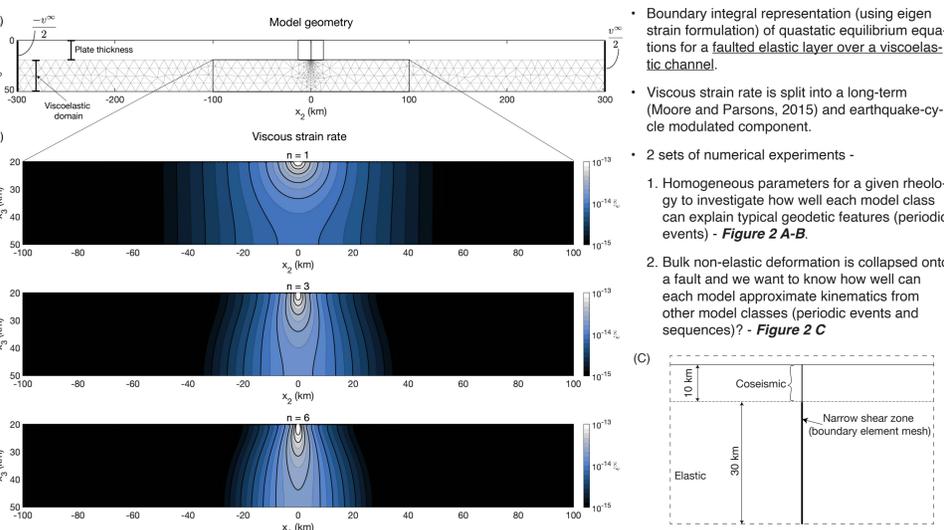


Figure 2. (A) Geometry of the first set of numerical experiments. (B) Long-term viscous strain rate as a function of n. (C) Model geometry for the second set of numerical experiments

- Boundary integral representation (using eigen strain formulation) of quasatic equilibrium equations for a faulted elastic layer over a viscoelastic channel.
- Viscous strain rate is split into a long-term (Moore and Parsons, 2015) and earthquake-cycle modulated component.
- 2 sets of numerical experiments -
 1. Homogeneous parameters for a given rheology to investigate how well each model class can explain typical geodetic features (periodic events) - Figure 2 A-B.
 2. Bulk non-elastic deformation is collapsed onto a fault and we want to know how well can each model approximate kinematics from other model classes (periodic events and sequences)? - Figure 2 C

Explaining interseismic & postseismic geodetic features for periodic cycles

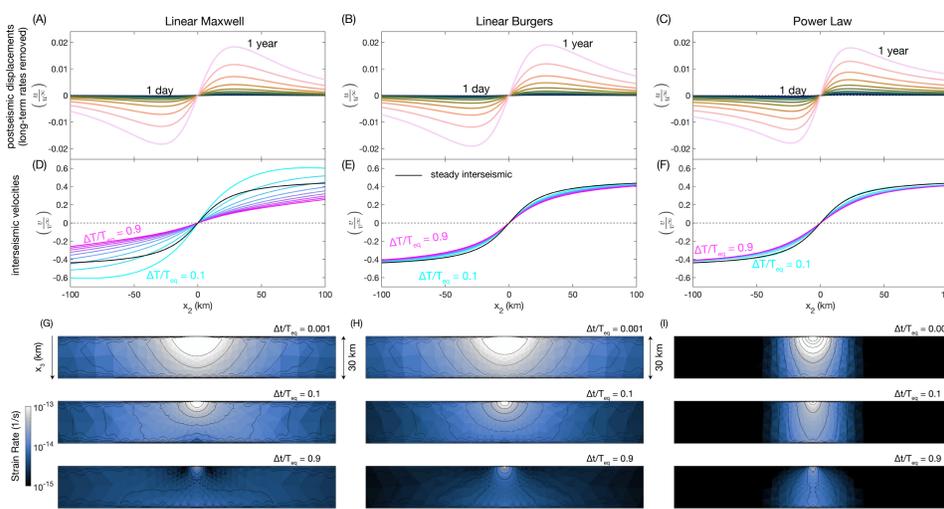


Figure 3. Surface predictions of postseismic displacements and interseismic velocities for different rheological models for a periodic earthquake cycle. The rheological parameters are chosen such that the cumulative postseismic after 1 year is nearly identical for all three models. (A)-(C) Cumulative postseismic displacements normalized by the coseismic slip amount $u_{eq} = v_{eq} T_{eq}$ for times varying from 1 day to 1 year. (D)-(F) Interseismic velocities compared to the steady interseismic expectation (black line). (G)-(I) Snapshots of internal viscous strain rates at different times during the earthquake cycle for each rheological model.

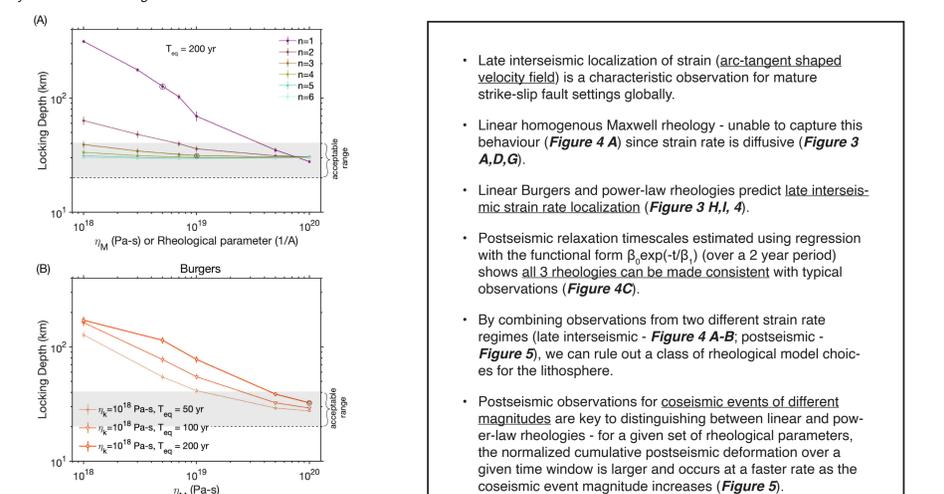


Figure 4. Compilation of late interseismic locking depths for various rheological choices. Locking Depth for (A) Linear Maxwell and power-law materials with n varying from 1 to 6 for $T_{eq} = 200$ years. (B) Locking depths for a linear Burgers rheology for a constant η_0 and varying η_1 and T_{eq} . Late interseismic locking depths show no dependence on η_1 . (C) Post-seismic relaxation times estimated over a 2-year period following the earthquake.

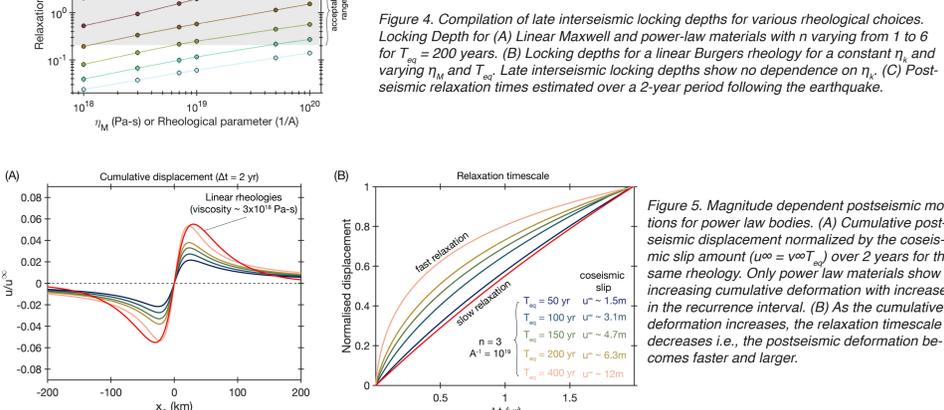
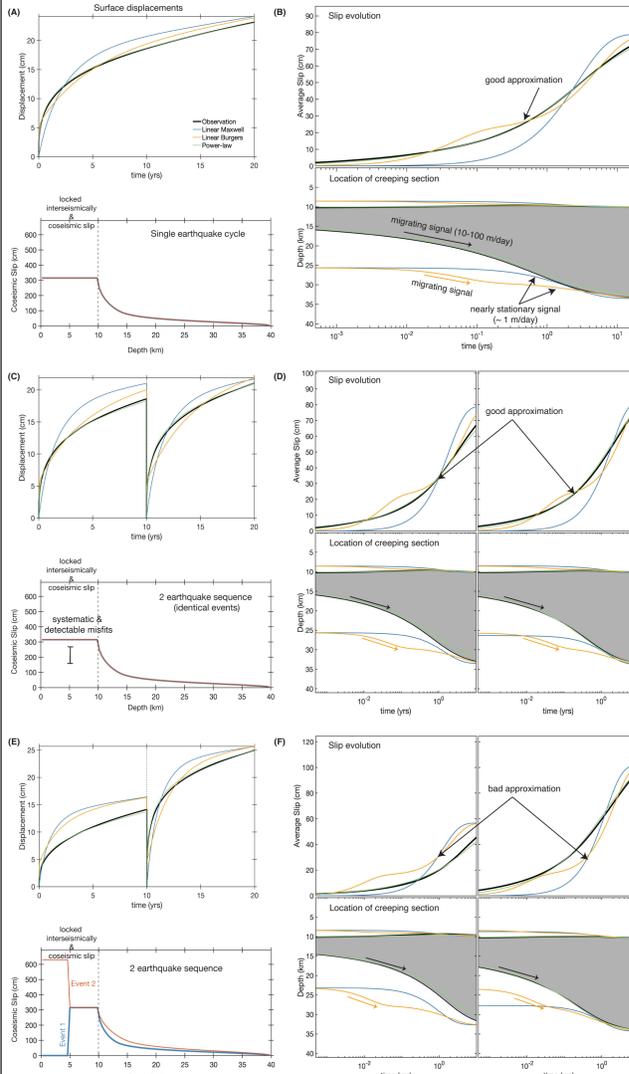


Figure 5. Magnitude dependent postseismic motions for power law bodies. (A) Cumulative post-seismic displacement normalized by the coseismic slip amount ($u_{eq} = v_{eq} T_{eq}$) over 2 years for the same rheology. Only power law materials show increasing cumulative deformation with increase in the recurrence interval. (B) As the cumulative deformation increases, the relaxation timescale decreases i.e., the postseismic deformation becomes faster and larger.

Observations from earthquake sequences can constrain the average rheological model of the lithosphere uniquely



Implications for time-dependent regional stress transfer

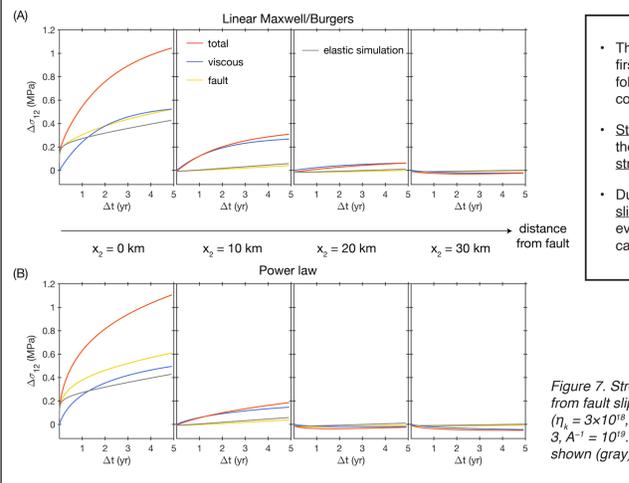


Figure 7. Stress change and decomposition into contributions from fault slip and viscous shear for (A) linear Burgers rheology ($\eta_1 = 3 \times 10^{18}$ Pa-s, $\eta_0 = 10^{20}$ Pa-s), and (B) power law rheology with $n = 3$, $A^{-1} = 10^{19}$. Total stress evolution from an elastic model is also shown (gray).

- For periodic cycles, the kinematics of heterogeneous linear Maxwell, homogeneous linear Burgers and power-law models are indistinguishable at the level of geodetic position uncertainty ~ 1 cm (Figure 6 A).
- Similar kinematics but dramatically different dynamics - location of peak creep rate and stress (Figure 6 B).
- Migrating stress front may drive after-shock activity (Figure 6 B). Could be an additional way to constrain parameters for a given rheology.
- Earthquake sequences where the coseismic events are of similar magnitudes (perturbation of pre-seismic conditions) may be insufficient to distinguish between linear Burgers and power-law rheologies (Figure 6 C,D).
- Events that vary by an order of magnitude allow clear distinction between various rheologies (Figure 6 E,F).